# The Arthur trace formula and spectral theory on locally symmetric spaces 

Werner Müller<br>University of Bonn<br>Institute of Mathematics

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## Introduction

The Selberg trace formula establishes a close relation between spectral and geometric data for finite volume locally symmetric spaces of rank 1 .
For a general reductive group $G$ over a number field $F$, Arthur, driven by Langlands' functoriallity conjectures, developed a trace formula for adelic quotients $G(F) \backslash G(\mathbb{A})$.
The key issue in Arthur's work is the comparison of the trace formulas of two different groups. However, it can also be used to study spectral problems on a single space. Such applications lead to new analytic problems related to the trace formula itself.

## 1. The Selberg trace formula

- G semisimple real Lie group with finite center of non-compact type
- $K \subset G$ maximal compact subgroup
- 「 $\subset G$ lattice
- $R_{\Gamma}$ right regular representation of $G$ in $L^{2}(\Gamma \backslash G)$, defined by

$$
\left(R_{\Gamma}(g) f\right)\left(g^{\prime}\right)=f\left(g^{\prime} g\right), \quad f \in L^{2}(\Gamma \backslash G)
$$

Main goal: Study of the spectral resolution of $\left(R_{\ulcorner }, L^{2}(\Gamma \backslash G)\right)$.
a) $\Gamma$ uniform lattice

Gelfand, Graev, Piateski-Shapiro: $R_{\Gamma}$ decomposes discretely

$$
R_{\Gamma}=\bigoplus_{\pi \in \widehat{G}} m_{\Gamma}(\pi) \pi .
$$

Let $f \in C_{c}^{\infty}(G)$. Define

$$
R_{\Gamma}(f)=\int_{G} f(g) R_{\Gamma}(g) d g
$$

Then $R_{\Gamma}(f)$ is an integral operator

$$
\left(R_{\Gamma}(f) \varphi\right)(g)=\int_{\Gamma \backslash G} K_{f}\left(g, g^{\prime}\right) \varphi\left(g^{\prime}\right) d g^{\prime}, \quad \varphi \in L^{2}(\Gamma \backslash G),
$$

with kernel

$$
K_{f}\left(g, g^{\prime}\right)=\sum_{\gamma \in \Gamma} f\left(g^{-1} \gamma g^{\prime}\right)
$$

Since $\Gamma \backslash G$ is compact, $R_{\Gamma}(f)$ is a trace class operator and

$$
\operatorname{Tr} R_{\Gamma}(f)=\int_{\Gamma \backslash G} K_{f}(g, g) d g=\int_{\Gamma \backslash G} \sum_{\gamma \in \Gamma} f\left(g^{-1} \gamma g\right) d g
$$

- break the sum over $\gamma$ into conjugacy classes $\{\gamma\}$ of $\Gamma$.

Let $\Gamma_{\gamma}$ and $G_{\gamma}$ be the centralizer of $\gamma$ in $\Gamma$ and $G$, respectively. The contribution of a conjugacy class $\{\gamma\}$ is

$$
\int_{\Gamma_{\gamma} \backslash G} f\left(g^{-1} \gamma g\right) d \dot{g}=\operatorname{vol}\left(\Gamma_{\gamma} \backslash G_{\gamma}\right) I(\gamma, f)
$$

where $I(\gamma, f)$ is the orbital integral

$$
I(\gamma, f)=\int_{G_{\gamma} \backslash G} f\left(g^{-1} \gamma g\right) d \dot{g}, \quad f \in C_{c}^{\infty}(G)
$$

Thus we get

$$
\operatorname{Tr} R_{\Gamma}(f)=\sum_{\{\gamma\}} \operatorname{vol}\left(\Gamma_{\gamma} \backslash G_{\gamma}\right) /(\gamma, f)
$$

On the other hand, by the result of Gelfand, Graev, and Piatetski-Shapiro, we get

$$
\operatorname{Tr} R_{\Gamma}(f)=\sum_{\pi \in \widehat{G}} m_{\Gamma}(\pi) \operatorname{Tr} \pi(f)
$$

Comparing the two expressions, we obtain
Trace formula (1. version):

$$
\sum_{\pi \in \widehat{G}} m_{\Gamma}(\pi) \operatorname{Tr} \pi(f)=\sum_{\{\gamma\}} \operatorname{vol}\left(\Gamma_{\gamma} \backslash G_{\gamma}\right) I(\gamma, f) .
$$

spectral side $=$ geometric side

- $I(\gamma, f)$ and $\operatorname{Tr} \pi(f)$ are invariant distributions on $G$, i.e., invarinat under $f \rightarrow f^{g}$, where $f^{g}\left(g^{\prime}\right)=f\left(g^{-1} g^{\prime} g\right)$.
- Fourier inversion formula can be used to express $I(\gamma, f)$ in terms of characters.


## The rank one case.

To make the trace formula useful, one has to understand the distributions $I(\gamma, f)$ and $\operatorname{Tr} \pi(f)$ and to express them in differential geometric terms. This is possible if the $\mathbb{R}$-rank of $G$ is 1 .

We specialize to: $G=\operatorname{SL}(2, \mathbb{R}), K=\mathrm{SO}(2)$.

- $\mathbb{H}=G / K$ upper half-plane, $\Gamma \subset G$ co-compact.

Let

$$
f \in C_{c}^{\infty}(G / / K)=\left\{f \in C_{c}^{\infty}(G): f\left(k_{1} g k_{2}\right)=f(g), k_{1}, k_{2} \in K\right\} .
$$

Then $\operatorname{Tr} \pi(f)=0$, unless $\pi$ has a $K$-fixed vextor. Hence

$$
\operatorname{Tr} \pi(f) \neq 0 \Leftrightarrow \exists s \in i \mathbb{R} \cap[-1,1]: \pi=\pi_{s}
$$

Let

$$
\Delta=-y^{2}\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}\right), \quad x+i y \in \mathbb{H}
$$

- hyperbolic Laplaceoperator on $\mathbb{H}$.
- $\Delta$ has discrete spectrum in $L^{2}(\Gamma \backslash \mathbb{H})$.

$$
\sigma(\Delta): 0=\lambda_{0}<\lambda_{1}<\lambda_{2}<\cdots \rightarrow \infty .
$$

- $m\left(\lambda_{j}\right)$ multiplicity of $\lambda_{j}$.

Frobenius reciprocity: $m\left(\pi_{s}\right)=m\left(\left(1-s^{2}\right) / 4\right)$ for $s \in i \mathbb{R} \cup[-1,1]$.
Let

$$
\mathcal{A}(f)(t)=\int_{\mathbb{R}} f\left(\left(\begin{array}{cc}
e^{t / 2} & x \\
0 & e^{-t / 2}
\end{array}\right)\right) d x
$$

be the Abel transform of $f$. Then $\mathcal{A}$ defines an isomorphism

$$
\mathcal{A}: C_{c}^{\infty}(G / / K) \rightarrow C_{c}^{\infty}(\mathbb{R})^{\text {even }}
$$

Moreover $\mathcal{A}(f)$ is closely related to the orbital interal of $f$ :

$$
I\left(a_{t}, f\right)=\frac{1}{\left|e^{t}-e^{-t}\right|} \mathcal{A}(f)(2 t), \quad a_{t}=\left(\begin{array}{cc}
e^{t} & 0 \\
0 & e^{-t}
\end{array}\right) .
$$

Let $h=\mathcal{A}(f)$. Then

$$
\widehat{h}(r)=\int_{G} f(g) \phi_{1 / 2+i r}(g) d g
$$

is the spherical Fourier transform of $f$, where $\phi_{\lambda}$ is the spherical function.
$f$ can be recovered from $h$ by Plancherel inversion:

$$
f(e)=\int_{\mathbb{R}} \widehat{h}(r) r \tanh (r) d r
$$

Moreover, using the polar decomposition $G=K A K$, it follows that

$$
\widehat{h}(r)=\operatorname{Tr} \pi_{2 i r}(f)
$$

Assumption: 「 torsion free

- $\gamma \in \Gamma-\{e\}$ is hyperbolic,
- $\{\gamma\}$ corresponds to unique closed geodesic $\tau_{\gamma}$ in $\Gamma \backslash \mathbb{H}$.
- $\ell(\gamma)=$ length $\left(\tau_{\gamma}\right)$.

Let $\gamma \sim\left(\begin{array}{cc}e^{t / 2} & 0 \\ 0 & e^{-t / 2}\end{array}\right)$. Then $\ell(\gamma)=t$.
Write the eigenvalues of $\Delta$ as

$$
\lambda_{j}=\frac{1}{4}+r_{j}^{2}, \quad r_{j} \in \mathbb{R} \cap i[-1 / 2,1 / 2] .
$$

Each $\gamma$ can be uniquely written as $\gamma=\gamma_{0}^{k}, k \in \mathbb{N}$, where $\gamma_{0}$ is primitive. Then

$$
\operatorname{vol}\left(\Gamma_{\gamma} \backslash G_{\gamma}\right)=\ell\left(\gamma_{0}\right)
$$

Selberg's trace formula (K-invariant form):

$$
\begin{aligned}
\sum_{j} m\left(\lambda_{j}\right) \widehat{h}\left(r_{j}\right)= & \frac{\operatorname{Area}(\Gamma \backslash \mathbb{H})}{2 \pi} \int_{\mathbb{R}} \widehat{h}(r) r \tanh (\pi r) d r \\
& +\sum_{\{\gamma\} \neq e} \frac{\ell\left(\gamma_{0}\right)}{e^{\ell(\gamma) / 2}-e^{-\ell(\gamma) / 2}} h(\ell(\gamma))
\end{aligned}
$$

- The kernel function $f \in C_{c}^{\infty}(G / / K)$ has been eliminated from the formula.
- $h \in C_{c}^{\infty}(\mathbb{R})$.
b)「 non-uniform

We assume that $\operatorname{vol}(\Gamma \backslash G)<\infty$ and $\Gamma \backslash G$ non-compact.

- $R_{\Gamma}(f)$ is not trace class
- $R_{\Gamma}$ does not decompose discretely.

Langlands's theory of Eisenstein series provides a decomposition into invariant subspaces

$$
L^{2}(\Gamma \backslash G)=L_{d}^{2}(\Gamma \backslash G) \oplus L_{a c}^{2}(\Gamma \backslash G)
$$

where

$$
R_{\Gamma}^{d}=\bigoplus_{\pi \in \widehat{G}} m_{\Gamma}(\pi) \pi
$$

and $L_{d}^{2}(\Gamma \backslash G)$ is the maximal invariant subspace, in which $R_{\Gamma}$ decomposes discretely.

- $L_{a c}^{2}(\Gamma \backslash G)$ is described in terms of Eisenstein series.

Theorem. (Ji, Mü, '98): For each $f \in C_{c}^{\infty}(G), R_{\Gamma}^{d}(f)$ is a trace class operator.

Therefore

$$
\operatorname{Tr} R_{\Gamma}^{d}(f)=\sum_{\pi \in \widehat{G}} m_{\Gamma}(\pi) \operatorname{Tr} \pi(f) .
$$

- In higher rank, there is no trace formula within this framework.

The rank one case: $G=S L(2, \mathbb{R}), \Gamma \subset G$ a non-uniform lattice.

- $\Delta$ has continuous spectrum: $[1 / 4, \infty)$,
- possible eigenvalues of $\Delta: 0=\lambda_{0}<\lambda_{1}<\cdots$,
- the only obvious eigenfunction is the constant function for which $\lambda=0$.
- continuous spectrum is described by Eisenstein series.


A hyperbolic surface with 3 cusps.

The surface $X$ can be compactified by adding $m$ points $a_{1}, \ldots, a_{m}$ :

$$
\bar{X}=X \cup\left\{a_{1}, \ldots, a_{m}\right\}
$$

- $\bar{X}$ is a closed Riemann surface.
- The points $a_{1}, \ldots, a_{m}$ are called cusps. They correspond to parabolic fixed points $p_{1}, \ldots, p_{m} \in \mathbb{R} \cup\{\infty\}$ of $\Gamma$.
- $a_{k} \mapsto E_{k}(z, s)$, Eisenstein series attached to $a_{k}$.

Example: $\Gamma=\operatorname{SL}(2, \mathbb{Z})$.

- $\Gamma \backslash \mathbb{H}$ has a single cusp $\infty$.
- Eisenstein series attached to $\infty$ :

$$
E(z, s)=\sum_{\gamma \in \Gamma_{\infty} \backslash \Gamma} \operatorname{Im}(\gamma z)^{s}=\sum_{(m, n)=1} \frac{y^{s}}{|m z+n|^{2 s}}, \quad \operatorname{Re}(s)>1
$$

## Properties:

- $E(\gamma z, s)=E(z, s), \gamma \in \operatorname{SL}(2, \mathbb{Z})$.
- $E(z, s)$ admits a meromorphic extension to $s \in \mathbb{C}$,
- $E(z, s)$ is holomorphic on $\operatorname{Re}(s)=1 / 2$,
- $\Delta E(z, s)=s(1-s) E(z, s)$.

It follows that $r \in \mathbb{R} \mapsto E(z, 1 / 2+i r)$ is a generalized eigenfunction.

## Scattering matrix

Fourier expansion of $E(z, s)$ :

$$
E(x+i y, s)=y^{s}+C(s) y^{1-s}+O\left(e^{-c y}\right)
$$

as $y \rightarrow \infty$. Sommerfeld radiation condition

- $y^{1 / 2+i r}$ incoming plane wave, $y^{1 / 2-i r}$ outgoing plane wave, $E(z, 1 / 2+i r)$ the distorted plane wave.
- $S(r)=C(1 / 2+i r)$ scattering matrix,
- $C(s)$ analytic continuation of the scattering matrix,

General case: $E_{k}(z, s), k=1, \ldots, m$, Eisenstein series. Fourier expansion of $E_{k}(z, s)$ in the cusp $a_{l}$ gives scattering matrix:

$$
C(s)=\left(C_{k l}(s)\right)_{k, l=1}^{m} .
$$

Let $\phi(s)=\operatorname{det} C(s)$.

$$
\frac{1}{4 \pi} \int_{\mathbb{R}} \widehat{h}(r) \frac{\phi^{\prime}}{\phi}(1 / 2+i r) d r
$$

contribution of the Eisenstein series to the trace formula.

- 「 has now parabolic elements

Parabolic contribution:

$$
\int_{\mathbb{R}} \widehat{h}(r) \frac{\Gamma^{\prime}}{\Gamma}(1+i r) d r .
$$

Selberg trace formula for non-uniform lattices:

$$
\begin{aligned}
& \sum_{j} \widehat{h}\left(r_{j}\right)-\frac{1}{4 \pi} \int_{-\infty}^{\infty} \widehat{h}(r) \frac{\phi^{\prime}}{\phi}(1 / 2+i r) d r+\frac{1}{4} \phi(1 / 2) h(0) \\
& =\frac{\operatorname{Area}(\Gamma \backslash \mathbb{H})}{4 \pi} \int_{\mathbb{R}} \widehat{h}(r) r \tanh (\pi r) d r+\sum_{\{\gamma\} \neq e} \frac{\ell\left(\gamma_{0}\right)}{e^{\ell(\gamma) / 2}-e^{-\ell(\gamma) / 2}} h(\ell(\gamma)) \\
& -\frac{m}{2 \pi} \int_{-\infty}^{\infty} \widehat{h}(r) \frac{\Gamma^{\prime}}{\Gamma}(1+i r) d r+\frac{m}{4} \widehat{h}(0)-m \ln 2 h(0)
\end{aligned}
$$

- Can be understood as relative trace formula


## II. Applications of the trace formula

## 1) Weyl's law and the existence of cups forms

Rank one case: $G=S L(2, \mathbb{R}), \Gamma \subset G$ non-uniform lattice. Let $0=\lambda_{0}<\lambda_{1} \leq \cdots$ be the eigenvalues of $\Delta, C(s)$ scattering matrix, $\phi(s)=\operatorname{det} C(s)$. Put

$$
N_{\Gamma}(\lambda)=\#\left\{j: \lambda_{j} \leq \lambda^{2}\right\}, \quad M_{\Gamma}(\lambda)=-\frac{1}{4 \pi} \int_{-\lambda}^{\lambda} \frac{\phi^{\prime}}{\phi}(1 / 2+i r) d r .
$$

Theorem 1 (Selberg): As $\lambda \rightarrow \infty$, we have

$$
N_{\Gamma}(\lambda)+M_{\Gamma}(\lambda)=\frac{\operatorname{Area}(\Gamma \backslash \mathbb{H})}{4 \pi} \lambda^{2}+O(\lambda \log \lambda)
$$

proof: (without remainder term)

- $k_{t}$ kernel of the heat operator $e^{-t \widetilde{\Delta}}$ on $\mathbb{H}$.
- $k_{t} \in \mathcal{C}^{1}(G / / K)$ (bi-K-invariant, integrable, rapidely decraesing functions).
- Selberg trace formula can be applied to $k_{t}$.

Let $h_{t}=\mathcal{A}\left(k_{t}\right)$ be the Abel transform. Then

$$
h_{t}(x)=\frac{1}{\sqrt{4 \pi t}} e^{-t / 4-x^{2} /(4 t)}, \quad \widehat{h}_{t}(r)=e^{-\left(1 / 4+r^{2}\right) t}
$$

If we insert $h_{t}$ in the trace formula, we get

$$
\sum_{j} e^{-t \lambda_{j}}-\frac{1}{4 \pi} \int_{\mathbb{R}} e^{-\left(1 / 4+r^{2}\right) t} \frac{\phi^{\prime}}{\phi}\left(\frac{1}{2}+i r\right) d r \sim \frac{\operatorname{Area}(X)}{4 \pi} t^{-1}
$$

as $t \rightarrow 0+$.

- For $\lambda \gg 0$, the winding number $M_{\Gamma}(\lambda)$ is monotonic increasing.
- Tauberian theroem $\Rightarrow$ Theorem.

A more sophisticated use of the trace formula gives an estimation of the remainder term.
First step is to estimate the number of eigenvalues in an interval. Hörmander's method.

## The scattering matrix for arithmetic groups

- In general, $N_{\Gamma}(\lambda)$ and $M_{\Gamma}(\lambda)$ can not be separated.
- For the principal congruence subgroup $\Gamma(N)$, the entries of the scattering matrix can be expressed in terms of known functions of analytic number theory.

Huxley: For $\Gamma(N)$ we have

$$
\phi(s)=(-1)^{\prime} A^{1-2 s}\left(\frac{\boldsymbol{\Gamma}(1-s)}{\boldsymbol{\Gamma}(s)}\right)^{k} \prod_{\chi} \frac{L(2-2 s, \bar{\chi})}{L(2 s, \chi)}
$$

where $k, l \in \mathbb{Z}, A>0, \chi$ Dirichlet character $\bmod k, k \mid N, L(s, \chi)$ Dirichlet $L$-function with character $\chi$.
Especially, for $N=1$ we have

$$
\phi(s)=\sqrt{\pi} \frac{\boldsymbol{\Gamma}(s-1 / 2) \zeta(2 s-1)}{\boldsymbol{\Gamma}(s) \zeta(2 s)},
$$

where $\zeta(s)$ denotes the Riemann zeta function.

Thus for $\Gamma(N)$ we get

$$
\left|\frac{\phi^{\prime}}{\phi}(1 / 2+i r)\right| \ll \log ^{k}(|r|+1), \quad r \in \mathbb{R}
$$

and therefore

$$
M_{\Gamma(N)}(\lambda)=O(\lambda \log \lambda)
$$

Theorem 2 (Selberg, 1956):

$$
N_{\Gamma(N)}(\lambda)=\frac{\operatorname{Area}(\Gamma(N) \backslash \mathbb{H})}{4 \pi} \lambda^{2}+O(\lambda \log \lambda), \quad \lambda \rightarrow \infty
$$

- For $\Gamma(N), L^{2}$-eigenfunctions of $\Delta$ with eigenvalue $\lambda \geq 1 / 4$ ( = Maass automorphic cusp forms) exist in abundance.
- For $\Gamma(1)=\mathrm{SL}(2, \mathbb{Z})$ no eigenfunction with eigenvalue $\lambda>0$ can be constructed explicitly.


## 2) Distribution of Hecke eigenvalues

$S_{k}(\Gamma(1))$ space of cusp forms of weight $k$.

$$
T_{n}: S_{k}(\Gamma(1)) \rightarrow S_{k}(\Gamma(1))
$$

the $n$-th Hecke operator.

- $S_{k}$ the set of all normalized Hecke eigenforms $f \in S_{k}(\Gamma(1))$.

Then

$$
T_{n} f=a_{f}(n) f, \quad f \in S_{k} .
$$

Put $\lambda_{f}(n)=n^{(1-k) / 2} a_{f}(n)$.
Deligne: $\quad \lambda_{f}(p) \in[-2,2]$ for $p$ prime.
Conjecture ( Serre): For each $h \in C([-2,2])$

$$
\frac{1}{\pi(x)} \sum_{p \leq x} h\left(\lambda_{f}(p)\right) \rightarrow \frac{1}{2 \pi} \int_{-2}^{2} h(t) \sqrt{4-t^{2}} d t, \quad x \rightarrow \infty
$$

Sato-Tate conjecture for modular forms.
Theorem (H. Nagoshi, 2006): Suppose that $k=k(x)$ satisfies $\frac{\log k}{\log x} \rightarrow \infty$ as $x \rightarrow \infty$. Then for every $h \in C([-2,2])$, we have

$$
\frac{1}{\pi(x) \# S_{k}} \sum_{\substack{p \leq x \\ f \in S_{k}}} h\left(\lambda_{f}(p)\right) \rightarrow \frac{1}{2 \pi} \int_{-2}^{2} h(t) \sqrt{4-t^{2}} d t, \quad x \rightarrow \infty
$$

4) Limit multiplicities
a) $\Gamma \subset G$ uniform lattice

- $\Gamma=\Gamma_{1} \supset \Gamma_{2} \supset \cdots \supset \Gamma_{n} \supset \cdots$ tower of normal subgroups of finte index, $\cap_{j} \Gamma_{j}=\{e\}$.

$$
R_{\Gamma_{j}}=\widehat{\bigoplus}_{\pi \in \widehat{G}} m\left(\Gamma_{j}, \pi\right) \pi
$$

- $S \subset \widehat{G}$ open, relative compact, regular for the Plancherel measure $\mu_{P L}$.

Put

$$
\mu_{j}(S)=\frac{1}{\operatorname{vol}\left(\Gamma_{j} \backslash G\right)} \sum_{\pi \in S} m\left(\Gamma_{j}, \pi\right)
$$

deGeorge-Wallach, Delorme: $\lim _{j \rightarrow \infty} \mu_{j}(S)=\mu_{P L}(S)$.
b) $\Gamma \subset G$ non-uniform lattice

Savin: $\pi \in \widehat{G}_{d}$.

$$
\lim _{j \rightarrow \infty} \mu_{j}(\{\pi\})=d(\pi)
$$

Clozel: weak version.

$$
\liminf _{j \rightarrow \infty} \mu_{j}(\{\pi\})>\varepsilon>0 .
$$

5) Low lying zeros of L-functions
$f \in S_{k}, L(s, f) L$-function attached to $f, \phi$ test function

$$
D(f, \phi)=\sum_{\gamma} \phi(\gamma)
$$

where $\gamma$ ranges over normalized zeros of $L(s, f) . \mathcal{F}(Q)$ family of $L$-functions depending on parameter $Q$.

$$
E(\mathcal{F}(Q), f)=\frac{1}{|\mathcal{F}(Q)|} \sum_{f \in \mathcal{F}(Q)} D(f, \phi)
$$

Behavior as $Q \rightarrow \infty$.
5) Jacquet-Langlands

Correspondence between automorphic forms of a quarternion algebra and $\mathrm{GL}(2)$.

$$
\begin{aligned}
(\Gamma \backslash \mathbb{H}) & \leftrightarrow\left(\Gamma^{\prime} \backslash \mathbb{H}\right) \\
\left\{\lambda_{j}, t_{p, j}\right\} & \leftrightarrow\left\{\lambda_{j}^{\prime}, t_{p, j}^{\prime}\right\}
\end{aligned}
$$

$\Gamma^{\prime} \backslash \mathbb{H}$ is a compact Riemann surface attached to a congruence quaternion group $\Gamma^{\prime}, \Gamma \backslash H$ non-compact congruence surface.

## III. Higher rank

- $S=G / K, \Delta$ Laplace operator on $\Gamma \backslash S$.
- $L_{\text {cus }}^{2}(\Gamma \backslash S) \subset L^{2}(\Gamma \backslash S)$ closure of the span of the space of cusp forms.
- $\Delta$ has discrete spectrum in $L_{\text {cus }}^{2}(\Gamma \backslash S)$.

$$
L_{\mathrm{dis}}^{2}(\Gamma \backslash S)=L_{\mathrm{cus}}^{2}(\Gamma \backslash S) \oplus L_{\mathrm{res}}^{2}(\Gamma \backslash S) .
$$

$N_{\Gamma}^{\text {cus }}(\lambda), N_{\Gamma}^{\text {res }}(\lambda)$ counting function of cuspidal and residual spectrum, resp.

## General results:

Theorem (Donnelly, '82): Let $d=\operatorname{dim} S$.

$$
\limsup _{\lambda \rightarrow \infty} \frac{N_{\Gamma}^{\text {cus }}(\lambda)}{\lambda^{d / 2}} \leq \frac{\operatorname{vol}(\Gamma \backslash S)}{(4 \pi)^{d / 2} \Gamma\left(\frac{d}{2}+1\right)}
$$

Theorem (Mü, '89): $N_{\Gamma}^{\text {res }}(\lambda) \ll \lambda^{2 d}, \quad \lambda \geq 1$.
Conjecture 1 (Sarnak): $\operatorname{rank}(S)>1$. Then $N_{\Gamma}^{\text {cus }}(\lambda)$ satisfies Weyl's law.
Conjecture 2: $N_{\Gamma}^{\text {res }}(\lambda) \ll \lambda^{(d-1) / 2}$.
Theorem (Lindenstrauss, Venkatesh): G split adjoint semisimple group over $\mathbb{Q}, G=\mathbf{G}(\mathbb{R}), \Gamma \subset \mathbf{G}(\mathbb{Q})$ a congrunece group, $d=\operatorname{dim} S$. Then

$$
N_{\Gamma}^{\text {cus }}(\lambda) \sim \frac{\operatorname{vol}(\Gamma \backslash S)}{(4 \pi)^{d / 2} \Gamma(d / 2+1)} \lambda^{d / 2}, \quad \lambda \rightarrow \infty
$$

- Confirms the conjectur of Sarnak in these cases.

Previous results:
S. Miller: $G=\operatorname{SL}(3, \mathbb{R}), \Gamma=\operatorname{SL}(3, \mathbb{Z})$,

Mü: $G=\operatorname{SL}(n, \mathbb{R}), \Gamma=\Gamma(N)$.

## Estimation of the remainder term

Theorem (Lapid, Mü, 2007): Let $S_{n}=\operatorname{SL}(n, \mathbb{R}) / \mathrm{SO}(n)$, $d=\operatorname{dim} S_{n}, N \geq 3$. Then
$N_{\Gamma(N)}^{\text {cus }}(\lambda)=\frac{\operatorname{vol}(\Gamma(N) \backslash S)}{(4 \pi)^{d / 2} \Gamma(d / 2+1)} \lambda^{d / 2}+O\left(\lambda^{(d-1) / 2}(\log \lambda)^{\max (n, 3)}\right)$.
Method: Combination of Hörmander's method and Arthur's trace formula.

Mœglin, Waldsburger, 1989: Description of the residual spectrum of $\mathrm{GL}(n)$.
Combined with Donnelly's estimate, we get
Theorem (Mœglin, Waldsburger, 1989): $S_{n}=\operatorname{SL}(n, \mathbb{R}) / \mathrm{SO}(n)$, $d=\operatorname{dim} S_{n}$.

$$
N_{\Gamma(N)}^{\mathrm{res}}(\lambda) \ll \lambda^{d / 2-1} .
$$

## Multidimensional version

- $G=$ NAK Iwasawa decomposition, $\mathfrak{a}=\operatorname{Lie}(A), H: G \rightarrow \mathfrak{a}$, $H($ nak $)=\log a, W=W(G, A)$.
- $\mathcal{D}(S)$ ring of invariant differential operators on $S$.

Harish-Chandra: $\mathcal{D}(S) \cong S\left(\mathfrak{a}_{\mathbb{C}}\right)^{W}$.
Thus, if

$$
\chi: \mathcal{D}(S)=S\left(\mathfrak{a}_{\mathbb{C}}^{*}\right)^{W} \rightarrow \mathbb{C}
$$

is a character. Then

$$
\chi=\chi_{\lambda} \leftrightarrow \lambda \in \mathfrak{a}_{\mathbb{C}}^{*} / W .
$$

For $\lambda \in \mathfrak{a}_{\mathbb{C}}^{*}$ let

$$
\mathcal{E}_{\mathrm{cus}}(\lambda)=\left\{\varphi \in L_{\mathrm{cus}}^{2}(\Gamma \backslash S): D \varphi=\chi_{\lambda}(D) \varphi, D \in \mathcal{D}(S)\right\}
$$

Lwt $m_{\text {cus }}(\lambda)=\operatorname{dim} \mathcal{E}_{\text {cus }}(\lambda)$. Then the cuspidal spectrum is defined as

$$
\Lambda_{\text {cus }}(\Gamma)=\left\{\lambda \in \mathfrak{a}_{\mathbb{C}}^{*} / W: m(\lambda)>0\right\} .
$$

- $\Lambda_{\text {cus }}(\Gamma) \cap i \mathfrak{a}^{*} / W$ is the tempered spectrum
- $\Lambda_{\text {cus }}(\Gamma)-\left(\Lambda_{\text {cus }}(\Gamma) \cap i \mathfrak{a}^{*} / W\right)$ the complementary spectrum.

Theorem (Lapid, Mü, 2007): Let $S_{n}=\operatorname{SL}(n, \mathbb{R}) / \mathrm{SO}(n)$ and $d_{n}=\operatorname{dim} S_{n}, \Omega \subset i \mathfrak{a}^{*}$ a bounded open subset with piecewise $C^{2}$ boundary, $\beta(\lambda)$ be the Plancherel measure. Then as $t \rightarrow \infty$

$$
\begin{array}{r}
\sum_{\lambda \in \Lambda_{\operatorname{cus}}(\Gamma(N)), \lambda \in t \Omega} m(\lambda)=\frac{\operatorname{vol}\left(\Gamma(N) \backslash S_{n}\right)}{|W|} \int_{t \Omega} \beta(\lambda) d \lambda \\
+O\left(t^{d_{n}-1}(\log t)^{\max (n, 3)}\right)
\end{array}
$$

and

$$
\sum_{\substack{\lambda \in \Lambda_{\operatorname{cus}}(\Gamma(N)) \\ \lambda \in B_{t}(0) \backslash i \mathfrak{a}^{*}}} m(\lambda)=O\left(t^{d_{n}-2}\right) .
$$

Duistermaat, Kolk, Varadarajan, 1979: This results holds for $G$ arbitrary, and $\Gamma \subset G$ a uniform lattice.

## IV. Problems

1) Generalize the results of Duistermaat-Kolk-Varadarajan on spectral asymptotics for compact locally symmetric spaces $\Gamma \backslash S$ to non-compact quotients where $\Gamma$ is a congruence subgroup. In particular, establish Weyl's law with a remainder term.
2) Analyze the spectral asymptotics for the Bochner-Laplace operator acting on the sections of a locally homogeneous vector bundle over $\Gamma \backslash S$ (i.e. automorphic forms with a given $K_{\infty}$-type).
3) Study the distribution of Hecke eigenvalues.
4) Study the distribution of low-lying zeros of L-functions of Hecke eigenforms of $\operatorname{SL}(n, \mathbb{Z}) \backslash \operatorname{SL}(n, \mathbb{R}) / \mathrm{SO}(n)$ with large eigenvalues.
5) Study the limiting behavior of the Laplace spectrum for towers
$\Gamma_{1} \supset \Gamma_{2} \supset \cdots$.

## V. The Arthur trace formula

- The Arthur trace formula is the main tool to study these problems in the higher rank case.
- General reductive group needs adelic framwork.
$G$ reductive algebraic group over $\mathbb{Q}, \mathbb{A}=\prod_{v}^{\prime} \mathbb{Q}_{v}$ ring of adels of
$\mathbb{Q}, G(\mathbb{A})=\prod_{v}^{\prime} G\left(\mathbb{Q}_{v}\right)$.
We study now the spectral resolution of the regular representation

$$
R: G(\mathbb{A}) \rightarrow \operatorname{Aut}\left(L^{2}(G(\mathbb{Q}) \backslash G(\mathbb{A}))\right)
$$

This is related to the previous framework as follows. Let $K_{f} \subset \prod_{p<\infty} G\left(\mathbb{Z}_{p}\right)$ be an open compact subgroup. Then

$$
G(\mathbb{Q}) \backslash G(\mathbb{A}) / K_{f} \cong \bigsqcup_{j}\left(\Gamma_{j} \backslash G(\mathbb{R})\right)
$$

$$
G(\mathbb{A})^{1}=\bigcap_{\chi \in X(G)_{\mathbb{Q}}} \operatorname{ker}|\chi|, \quad G(\mathbb{A})=G(\mathbb{A})^{1} \cdot A_{G}(\mathbb{R})^{0}
$$

The (non-invariant) trace formula is an identity of distributions on $G(\mathbb{A})^{1}$

$$
\begin{aligned}
\sum_{\chi \in \mathfrak{X}} J_{\chi}(f) & =\sum_{\mathfrak{o} \in \mathfrak{O}} J_{\mathfrak{0}}(f), \quad f \in C_{0}^{\infty}\left(G(\mathbb{A})^{1}\right) \\
\text { spectral side } & =\text { geometric side }
\end{aligned}
$$

- $\mathfrak{X}$ set of cuspidal data; equivalence classes of $(M, \rho), M$ Levi factor of rational parabolic subgroup, $\rho$ cuspidal automorphic representation of $M(\mathbb{A})^{1}$.
- $\mathfrak{O}$ set of equivalence classes in $G(\mathbb{Q}), \gamma \sim \gamma^{\prime}$, if $\gamma_{s}$ and $\gamma_{s}^{\prime}$ are $G(\mathbb{Q})$-conjugate.


## Spectral side

- $J_{\chi}$ is derived from the constant terms of Eisenstein series and generalizes

$$
\frac{1}{4 \pi} \int_{\mathbb{R}} \widehat{h}(r) \frac{\phi^{\prime}}{\phi}(1 / 2+i r) d r
$$

- $P \subset G \mathbb{Q}$-parabolic subgroup, $P=M_{P} N_{P}$ Levi decomposition
- $A_{P} \subset M_{P}$ split component of the center of $M_{P}, \mathfrak{a}_{P}=\operatorname{Lie}\left(A_{P}\right)$
- $\mathcal{A}^{2}(P)$ square integrable automorphic forms on $N_{P}(\mathbb{A}) M_{P}(\mathbb{Q}) \backslash G(\mathbb{A})$
- $Q=M_{Q} N_{Q} \mathbb{Q}$-parabolic subgroup of $G, M_{P}=M_{Q}=M$.

$$
M_{Q \mid P}(\lambda): \mathcal{A}^{2}(P) \rightarrow \mathcal{A}^{2}(Q), \quad \lambda \in \mathfrak{a}_{P, \mathbb{C}}^{*}
$$

intertwining operator, meromorphic function of $\lambda$, main ingredient of $J_{\chi}$.

- $\pi \in \Pi\left(M(\mathbb{A})^{1}\right)$ determines subspace $\mathcal{A}_{\pi}^{2}(P) \subset \mathcal{A}^{2}(P)$ of automorphic forms which transform according to $\pi$
- $M_{Q \mid P}(\lambda, \pi)$ restriction of $M_{Q \mid P}(\lambda)$ to $\mathcal{A}_{\pi}^{2}(P)$.
- $\rho_{\pi}(P, \lambda)$ induced representation of $G(\mathbb{A})$ in $\overline{\mathcal{A}}_{\pi}^{2}(P)$.

Let $P$ be maximal parabolic and $\bar{P}$ the opposite parabolic group. Then the following integral-series is part of the spectral side
$\sum_{\pi \in \Pi_{\operatorname{cus}}\left(M(\mathbb{A})^{1}\right)} \int_{-\infty}^{\infty} \operatorname{Tr}\left(M_{\bar{P} \mid P}(i \lambda, \pi)^{-1} \frac{d}{d z} M_{\bar{P} \mid P}(i \lambda, \pi) \rho_{\pi}(P, i \lambda, f)\right) d \lambda$
Problem: Absolute convergence of the integral-series.
$\pi=\otimes_{v} \pi_{v}, \phi \in \mathcal{A}_{\pi}^{2}(P), \phi=\otimes_{v} \phi_{v} . S$ finite set of places, containing $\infty$, such that $\phi_{v}$ is fixed under $G\left(\mathbb{Z}_{p}\right)$ for $p \notin S$.
There exist finite-dimensional representations $r_{1}, \ldots, r_{m}$ of ${ }^{L} M$ such that

$$
M_{\bar{P} \mid P}(s, \pi) \phi=\bigotimes_{v \in S} M_{\bar{P} \mid P}\left(s, \pi_{v}\right) \phi_{v} \otimes \bigotimes_{v \notin S} \tilde{\phi}_{v} \cdot \prod_{i=1}^{m} \frac{L_{S}\left(i s, \pi, \widetilde{r}_{i}\right)}{L_{S}\left(1+i s, \pi, \widetilde{r}_{i}\right)}
$$

where

$$
L_{S}(s, \pi, r)=\prod_{v \notin S} L\left(s, \pi_{v}, r_{v}\right), \quad \operatorname{Re}(s) \gg 0
$$

is the partial automorphic $L$-function attached to $\pi$ and $r$.

- This reduces the problem to the estimation of the number of zeros of $L_{s}\left(s, \pi, \widetilde{r}_{j}\right)$ in a circle of radius $T$ as $T \rightarrow \infty$.
- Need to control the constants in terms of $\pi$.

Lapid, Mü, 2008: In general, the study of the distribution $J_{\chi}$ can be reduced to the study of integrals as above associated to maximal parabolics in Levi subgroups.

Theorem (Lapid, Mü, 2008): For every reductive group G, the spectral side of the trace formula is absolutely convergent.
Mü, Speh, Lapid, 2004: $G=G L(n)$.

- This is a first step.
- The intended applications of the trace formula to spectral problems require a finer analysis of the $L$-functions.

For $\mathrm{GL}(n)$ the relevant $L$-functions are the Rankin-Selberg $L$-functions $L\left(s, \pi_{1} \times \pi_{2}\right)$ attached to cuspidal automorphic representations $\pi_{i}$ of $\operatorname{GL}\left(n_{i}, \mathbb{A}\right), i=1,2, n=n_{1}+n_{2}$.

Jacquet, Shahidi, Mœglin/Waldspurger, ...: completed L-function $\Lambda\left(s, \pi_{1} \times \pi_{2}\right)$ has at most simple poles at $s=0,1$, $s(1-s) \wedge\left(s, \pi_{1} \times \pi_{2}\right)$ is entire of order 1 , satisfies functional equation.

## Geometric side

The distributions $J_{0}$ are given in terms weighted orbital integrals. In general, they are difficult to define. A special case is

$$
\int_{G_{\gamma} \backslash G} f\left(g^{-1} \gamma g\right) w(g) d g
$$

where $w(g)$ is a certain weight function.

- Weighted orbital integrals are non-invariant distributions.

